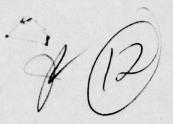
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THEORY OF POWER AUGMENTED RAM LIFT AT ZERO FORWARD SPEED

by

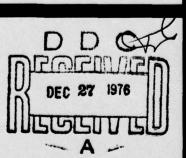
Roger W. Gallington and Harvey R. Chaplin

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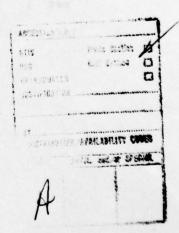
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but applying to larger ratios of propulsive jet area to leading edge air gap area (height of wing above ground times wing span) than were employed in the experiments. In this report, simple momentum theory is applied to obtain performance estimates for a range of jet area ratios from zero to the limiting cases of the potential flow solutions. Reasonable agreement with experiment is found, suggesting that two-dimensional flow models may give useful estimates of the performance potential of such systems. It is found that, under conditions such that a large fraction of the propulsion system thrust is recovered as useful propulsive thrust, the ratio of lift to thrust is very sensitive to the jet area ratio and optimum performance is obtained for jet area ratios approaching unity. Additional experiments are needed, covering the full range of jet area ratios, to evaluate the practical potential of power augmented ram lift systems.

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NOTATION

| Symbol | 2D Definition | Assumed 3D Equivalent |
|------------------|---|---|
| c | Wing chord | Wing area |
| h | Height of wing lower surface above ground | Area: (height of wing lower surface) X (wing span) |
| t ₁ | Thickness of incoming jet upstream of wing | Cross-sectional area of incoming jet upstream of wing |
| t ₂ | Thickness of departing jet downstream of wing | Cross-sectional area of depart- ing jet downstream of wing |
| β | Angle at which reentrant portion of jet leaves measured from original jet velocity vector | Angle at which reentrant portion of jet leaves measured from original jet velocity vector |
| Pc | Pressure exerted on wing lower surface in midchord region | Pressure exerted on wing lower surface in midchord region |
| q _j | Dynamic pressure of undisturbed jet | Dynamic pressure of undisturbed jet |
| C _p | Pressure coefficient: P _c /q _j | Pressure coefficient: P _c /q _j |
| т _ј | Jet thrust: 2qjt ₁ | Jet thrust |
| T _{net} | Net thrust: T minus drag force on wing | Net thrust: T minus drag force on wing |
| C _T | Thrust recovery coefficient: Tnet /Tj | Thrust recovery coefficient: Tnet Tj |
| L | Lift on wing | Lift on wing |

ABSTRACT

Recent experiments by Huffman and Jackson demonstrated the possibility of an air cushion operating mode for an aircraft with low aspect ratio end-plated wings wherein the propulsion system is mounted forward of the wing and exhausted into the space between the wing and the ground to generate high pressures for wing lift at zero forward speed. Gallington2 has derived two-dimensional potential flows relevant to this phenomenon, but applying to larger ratios of propulsive jet area to leading edge air gap area (height of wing above ground times wing span) than were employed in the experiments. In this report, simple momentum theory is applied to obtain performance estimates for a range of jet area ratios from zero to the limiting cases of the potential flow solutions. Reasonable agreement with experiment is found, suggesting that twodimensional flow models may give useful estimates of the performance potential of such systems. It is found that, under conditions such that a large fraction of the propulsion system thrust is recovered as useful propulsive thrust, the ratio of lift to thrust is very sensitive to the jet area ratio and optimum performance is obtained for jet area ratios approaching unity. Additional experiments are needed, covering the full range of jet area ratios, to evaluate the practical potential of power augmented ram lift systems.

ADMINISTRATIVE INFORMATION

This investigation was authorized and funded by the Naval Air Development Center under Project SSH15, Program Element 63534N, and Work Unit 1-1612-008.

INTRODUCTION

The experiments of Reference 1 demonstrate an interesting potential for generating wing lift at zero speed (and presumably at low speed) by directing the propulsion system exhaust into the space between the wing and the ground. For example, lift forces more than six times greater

¹Huffman, J. K. and C. M. Jackson, Jr., "Investigation of the Static Lift Capability of a Low-Aspect-Ratio Wing Operating in a Powered Ground Effect Mode," NASA TM X-3031 (Jul 1974).

²Gallington, R. W., "Sudden Deceleration of a Free Jet at the Entrance of a Channel," DTNSRDC Departmental Report ASED 350 (Jan 1976).

than the propulsion system thrust were generated at zero speed while simultaneously recovering about 60 percent of the propulsion system thrust as useful propulsive thrust.

This phenomenon might prove to be of great practical interest in future water-based aircraft design by providing means for avoiding high hydrodynamic drag during takeoff and high impact loads during takeoff and landing. Also, an air cushion cruise mode might be of interest to provide surface mobility over short ranges or under sea state conditions which prevented takeoff.

The available experimental data cover very limited ranges of the geometric parameters which appear to be important. It is desired to construct theoretical flow models which would provide insight into the phenomenon, estimates of performance limitations, and guidance for further experiments.

Two-dimensional potential flow solutions obtained in Reference 2 appear to be definitely relevant, and in fact may represent "optimum" cases of the power augmented ram lift phenomenon in certain important senses. However, they cover a range of jet area ratios which are not directly comparable to the experiments. Moreover, the experiments (and probably designs of practical interest as well) involve flows which are definitely <u>not</u> two-dimensional.

The step taken in this report will be to develop two-dimensional theoretical estimates which can be directly compared with the available experimental data in order to tentatively confirm that two-dimensional analyses may have a useable degree of applicability to practical cases.

MOMENTUM THEORY

In Reference 2, potential flow solutions were obtained for two classes of flows which will be referred to here as the "filled duct" class and the "overfilled duct" class, as illustrated in Figure 1.

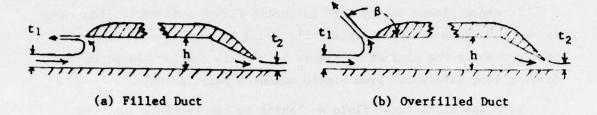


Figure 1 - Potential Flow Classes

These solutions have the property that the pressure recovery coefficient $C_p = 1 - (t_2/h)^2$ (which can be derived very simply from one-dimensional fluid mechanics principles). For a given value of C_p , the minimum admissible value of jet area ratio t_1/h corresponds to the filled-duct class. In other words, the filled-duct class is a limiting case of the overfilled duct class. To investigate cases of still smaller jet area ratios, we must adopt a modified flow model. The simplest, in the sense of involving fewest theoretical difficulties, is illustrated in Figure 2.

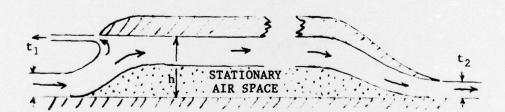


Figure 2 - Underfilled Duct Class

The incoming jet is assumed to curve away from the ground as it approaches the wing and to stagnate against the bottom of the wing near the leading edge splitting into one branch which reemerges forward parallel to the ground (as in the filled duct case) and a second which follows the bottom surface of the wing rearward until, near the trailing edge, it is led back into a tangent (nonstagnating) contact with the ground before reaccelerating to the original jet velocity at the trailing edge.

It seems highly likely that potential flow solutions of this class exist, but since it appears that they would be much more difficult to derive than in the classes of Figure 1, we will content ourselves with drawing some inferences from simple momentum relationships.

From one-dimensional fluid mechanics we can conclude that the thickness of the jet along the lower surface of the wing at midchord is $t_2/\sqrt{1-C_p}$ and its velocity is $V_j \sqrt{1-C_p}$. A momentum balance between a vertical plane upstream of the wing and one through the midchord thus takes the form

$$2q_{j}t_{1} + 2q_{j}(t_{1}-t_{2}) = P_{c}h + 2q_{j}t_{2}\sqrt{1-C_{p}}$$

$$C_{p} = 4t_{1}/h - 2(t_{2}/h(1+\sqrt{1-C_{p}}))$$
(1)

whereas this expression is rather awkward to solve for C_p directly, numerical solutions are readily extracted.

The thrust relationships are the same as in the filled-duct class, i.e.,

$$T_{j} = 2q_{j}t_{1}$$
 $T_{net} = 2q_{j}t_{2} - 2q_{j}(t_{1} - t_{2})$
 $C_{T} = 2t_{2}/t_{1} - 1$ (2)

There is a further difficulty in comparing these results to the experimental results of Reference 1 in that those experiments employed a compressed-air jet from a thin rectangular nozzle located many nozzle thicknesses upstream of the wing. The jet arriving at the wing is thus a fully-developed turbulent jet with a velocity distribution which is nonuniform and not precisely known. However, assuming momentum is conserved in the experimental jet, this problem can be

circumvented by working with force and momentum relationships rather than pressure relationships.

Noting that

$$L = q_j C_p c$$
 (assuming c >> h)

$$T_1 = 2q_1t_1$$

then

$$L/T_{j} = \frac{C_{p}}{2t_{1}/h} \frac{c}{h}$$

and the quantity

$$\frac{L}{T_i} \frac{h}{c} = \frac{C_p}{2t_1/h} \tag{3}$$

(which is readily evaluated numerically for given values of $C_{\rm T}$ and t_2/h from Equations (1) and (2)) lends itself to qualitative comparison with the available experimental data.

This comparison is made in Figure 3. It can be seen that, for the cases of moderate flap deflection, the agreement is reasonably good. Caution is advisable in drawing conclusions from this agreement since we are comparing uniform-jet, two-dimensional theory with distinctly nonuniform-jet, three-dimensional experiments. Also, the quantity h, which in the theory has been defined as distance from the ground to a parallel flat lower surface of the airfoil, is not directly definable for the experimental airfoil, which has a convex lower surface. (For plotting the experimental data, h was taken to be the height of the theoretical trailing edge of the airfoil above the ground.) Nevertheless, the two-dimensional theory seems to provide a reasonable basis for a tentative and cautious assessment of the potential of power augmented ram lift systems pending acquisition of a broader experimental data base.

Even if these uncertainties regarding applicability of two-dimensional analysis to practical cases are ultimately resolved, there are two further effects which clearly need to be taken into account: (1) viscous effects, expecially the additional thrust loss due to friction between the jet and the wing, and (2) effect of nonuniform pressure, expecially near the leading and trailing edges of the wing. (The theory assumes uniform pressure of magnitude ${\rm C}_{\rm pq}$ over the entire lower surface for purposes of computing lift, which is admissable only if h/c is very small.

The relationships available for two-dimensional analyses are summarized in Table 1.

CONCLUSIONS

- Theoretical relationships developed in Reference 2 for estimating the performance of power augmented ram lift systems have been extended to include all possible values of the jet area ratio.
- These relationships have been found to be in reasonable agreement with experimental results reported in Reference 1 (which are confined, however, to a limited range of rather small jet area ratios).
- While this agreement with experiment is hardly conclusive, it seems reasonable to use the theoretical relationships (with due caution and appropriate corrections) for preliminary studies of power-augmented ram lift systems, pending the availability of a more complete experimental data base.
 - Experiments to provide a more complete data base are needed.

TABLE 1 - FORMULAS FOR 2D ANALYSES

Section 8

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| | | | | | ~ | | | |
|------------------|-------------------|-------------------------------|---|---|--|---------------------------------------|---|-------------------------------|
| Underfilled Duct | 1 | 180° | independent $\frac{24t_2/h}{(1+t_2/h)^2}$ | | $1 - \left[t_2/h + \sqrt{(1+t_2/h)^2 - 4t_1/h}\right]^2$ | $2t_2/t_1-1$ | C _p /(2t ₁ /h) | |
| Overfilled Duct* | endent variable 5 | independent variable ≤ 0 | $2(1-\cos \beta) t_2/h$ 1- $(2\cos \beta) t_2/h + (t_2/h)^2$ | (t ₂ /h)/(t ₂ /t ₁) | /h)² | $(1-\cos \beta) t_2/t_1 + \cos \beta$ | $c_{\mathrm{p}}/(2\epsilon_{1}/\mathrm{h})$ | |
| Filled Duct* | indepe | 180° | $\frac{4t_2/h}{(1+t_2/h)^2}$ | | $1 - (t_2/h)^2$ | 2t ₂ /t ₁ -1 | $c_p/(2t_1/h)$ | $=\sqrt{2(1-C_{\mathrm{T}})}$ |
| Quantity | t ₂ /h | 80. | t ₂ /t ₁ | t,/h | ပ ^{ဇၗ} | $^{\rm C}_{ m T}$ | $\frac{L}{1}\frac{h}{c}$ | |

*From Reference 2.

 $^{^{\}text{h}}$ Note that the formulas assume h/c << 1. The effect of finite h/c is probably to reduce L/T $_{\rm J}$ somewhat.

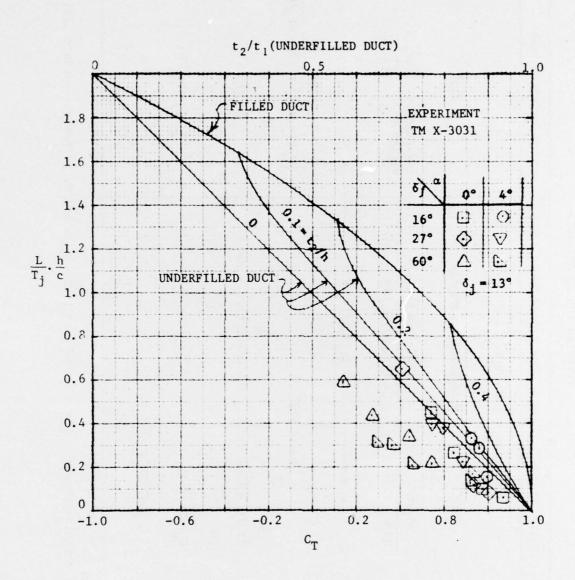


Figure 3 - Comparison with Experiment

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